

A new method to measure the shape of short THz pulses: Differential Electronic Gating

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Abstract - A simple experimental method has been developed to determine the shape of repetitive picosecond THz pulses in the presence of jitter in the trigger signal. This method, a modification of the recently reported Differential Optical Gating (DOG) method, is based on the femtosecond electronic gating of a high-frequency sequential oscilloscope. Preliminary tests have been performed on pulses from the free electron laser FELIX as well as on pulses from a mode locked p-Ge laser.

Introduction

The use of non-linear detection techniques to determine the temporal shape of picosecond THz pulses is in general not feasible due to the small optical susceptibilities of materials and/or low pulse intensities in this wavelength region. There do exist two *linear* techniques to analyze short THz pulses. These use high intensity femtosecond optical pulses to *create* as well as to *detect* - through optical or electronic gating - broadband THz pulses [1,2]. The intrinsic synchronization between THz and gating pulse results in a very good time resolution for those techniques. Other sources for the creation of short THz pulses, such as the free electron laser and the mode locked p-Ge laser, however, do not rely on optical pumping. The use of a gating technique, therefore, causes serious problems to synchronize the THz and gating pulses. A disadvantage of this DOG technique is that, in order to obtain a time resolution of the order of picoseconds, a mode-locked Ti:Sapphire laser system is needed. We have developed a simple technique, based on the DOG principle, that does not rely on optical, but on *electronic* gating. This work was started in order to determine the precise pulse width of our mode-locked p-Ge THz laser, which, until now was impossible because of the limited bandwidth of the 6 GHz single shot oscilloscope used. We show that, with this "Differential Electronic Gating" (DEG) method together with an ultra fast room temperature detector, a time resolution in the picosecond range is possible.

Principle of the differential technique

In the DEG technique the THz pulse is detected by one fast detector. The signal is *electronically* gated and recorded at two slightly different times t and $t + \Delta t$. In that way, both the average intensity $\bar{I} = [I(t+\Delta t)+I(t)]/2$ and the time derivative of the intensity $F(\bar{I}) = [I(t+\Delta t)-I(t)]/\Delta t$ are determined. For a series of THz pulses, this yields a set of derivative *versus* intensity

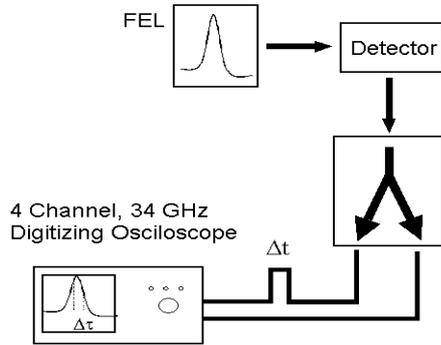


Fig. 1: Principal setup for the DEG method.

In order to obtain the pulse duration, one can take advantage of the fact that the density of data points indicates the relative amount of time spent in a given part of the $F-\bar{I}$ curve. By comparing the density in the peak region ($\bar{I} = \text{max.}$ and $F = 0$) with the density in the region, where $F \neq 0$, the time duration of the pulse peak can be inferred.

Experimental

The THz pulses are focused on a fast room temperature superlattice detector [3] consisting of a GaAs/AlAs superlattice mounted in a corner cube reflector with a long wire antenna. The video output of the detector is amplified with a 28 ps rise time amplifier and split by a Picosecond Pulse Labs. Model 5335 high frequency power splitter. One signal is delayed with respect to the other, using two coaxial lines with a slightly different length.(see fig. 1) The signals $I(t+\Delta t)$ and $I(t)$ are measured at the same time on two different channels of a HP 54120B gating oscilloscope with 54123A front-end, featuring a 34 GHz bandwidth and a 200 fs electronic gate width. Pulse shapes have also been studied using a single shot 6 GHz bandwidth Tektronix 7250 oscilloscope. The time delay between the channels, that determines the time scale of the reconstructed pulse, is measured by sending the scope's TDR pulse into the power splitter, and comparing the time delay of the signals on the two channels. At the same time the accuracy of the power splitting is checked. Using the normal sequential data collection routine of the oscilloscope, a set of data $\{I(t+\Delta t), I(t)\}$ is taken at each macropulse, with a 0.2 ps stepwise increase of the time delay between external trigger and electronic gate.

Results

A first test of the DEG technique has been performed using pulses at a wavelength of 150 μm from the free electron laser "FELIX", set at a 4 μs long macropulse and a micropulse repetition frequency of 1 GHz. For this experiment, the laser was slightly detuned to create a double pulse structure [4]. For this test $\Delta t=7$ ps was taken and the FELIX macropulse trigger was used which has a jitter of about 1 ns; nearly two orders of magnitude larger than the FWHM of the THz pulse . In the data from the first

data, $\{F(\bar{I})-\bar{I}\}$, which is *not* influenced by jitter in the trigger, as time as explicit variable has been eliminated. The actual pulse intensity as a function of time can now be reconstructed by integration: $\int_{\bar{I}_0}^{\bar{I}_1} d\bar{I} / F(\bar{I}) = t(\bar{I}_1) - t(\bar{I}_0)$. It must be noted that this integral is only well defined, if $F(\bar{I})$ does not contain any zeros. As at the pulse maximum the derivative will be zero, it is only possible to reconstruct rising and falling edges of a pulse shape.

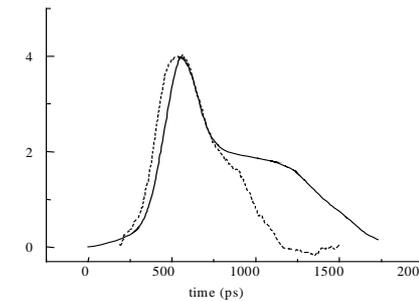
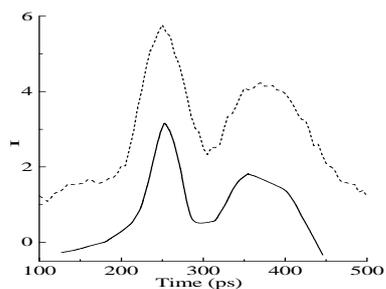
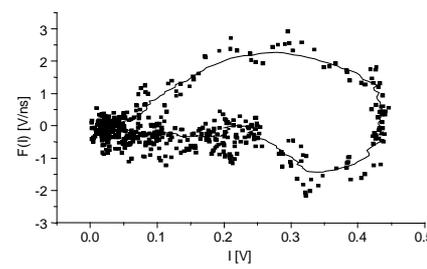
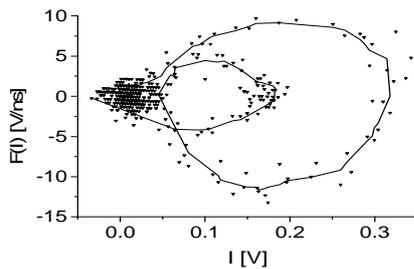
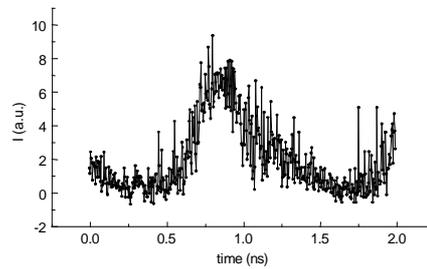
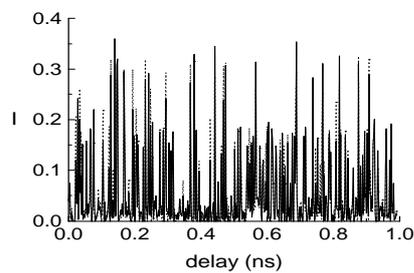


Fig. 2: FELIX double pulse:

- a: Sequential data with macropulse trigger; two channels with $\Delta t = 7$ ps.
- b: F versus I plot.
- c: Reconstructed pulse with 40 ps FWHM + 6 GHz scope signal (dotted line).

Fig. 3: Mode-locked p-Ge laser output:

- a: Sequential data with optical trigger.
- b: F versus I plot with $\Delta t = 40$ ps.
- c: Reconstructed pulse (solid line) and averaged sequential signal (dotted line).

channel given in fig. 2a., therefore, no pulse shape can be distinguished at all. In Fig. 2b the full data set $\{F-I\}$ plot is shown; a double loop curve is seen, characteristic for a pulse structure with two unequal pulses. The reconstructed double pulse is given in Fig. 2c. Because of the finite overlap between the two pulses, no zero intensity region in between the pulses occurs, and thus the distance between the pulses can be unambiguously determined from the experimental data. Also shown in this last figure is the pulse shape as determined with the single shot 6 GHz oscilloscope. The FWHM of the largest reconstructed pulse is about 40 ps, mainly set by the amplifier rise time. The FWHM value as observed with the 6 GHz scope is 56 ps; slightly larger due to the 50 ps rise time of that device. The second experiment was performed using the 175 μm

output of our mode-locked p-Ge THz laser [5]. Earlier experiments using the 6 GHz oscilloscope proved that that source was able to produce 60 ps FWHM pulses. In view of the 50 ps rise time of that detection system, the actual width was expected to be considerably shorter. In this experiment we used a 40 ps time delay between the two channels, and the optical signal itself was used as a trigger. In fig. 3a. the sequential data from two channels are shown, whereas in fig. 3b. the F versus I plot of a large data set is given. The reconstructed shape is shown in fig. 3c, together with a pulse shape obtained from an averaged sequential scan. Although the FWHM value of the reconstructed pulse is slightly smaller, the overall shape is worse than that obtained from the direct scan. Clearly this method does not work properly in this case. The reason for that is the strong shot to shot variation of both the intensity and the shape of the pulse resulting from strong mode beating effects of the laser.

Conclusions

We have shown that the use of the differential electronic gating (DEG) method enables the detection of repetitive picosecond THz pulses, even when the jitter of the synchronization signal is much larger than the width of the pulse. The time resolution of this method is limited by the electronic rise time of optical detector and electronics. The shot to shot reproducibility of the pulse shape, however, has to be good - like in all sequential recording techniques - in order to obtain an optimal result.

Acknowledgments

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